

# Probing QCD with top-quark pairs at CMS

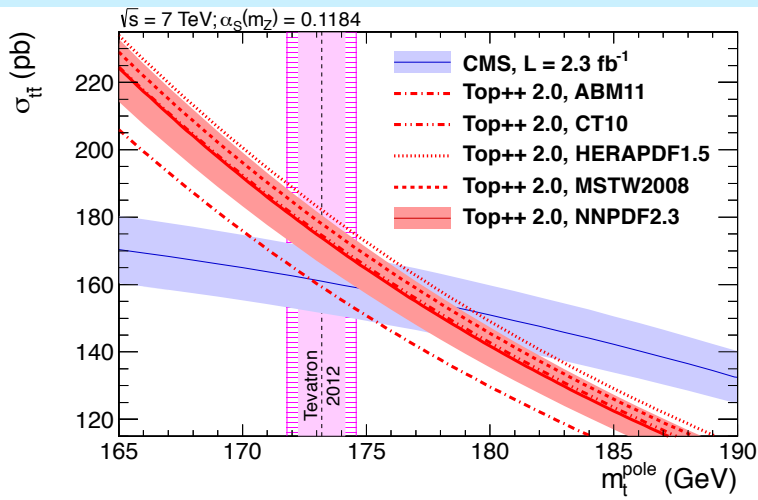
determination of  $m_t$  and  $\alpha_s$  using top-pair cross section

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The inclusive cross section for top-pair production as measured by the CMS experiment in  $pp$  collisions at  $\sqrt{s} = 7$  TeV is compared to the QCD prediction at NNLO using various PDF sets. For each PDF set, the pole mass of the top quark,  $m_t$ , and the strong coupling constant,  $\alpha_s(M_Z)$ , are extracted. This is the first determination of  $\alpha_s(M_Z)$  from top-quark production.

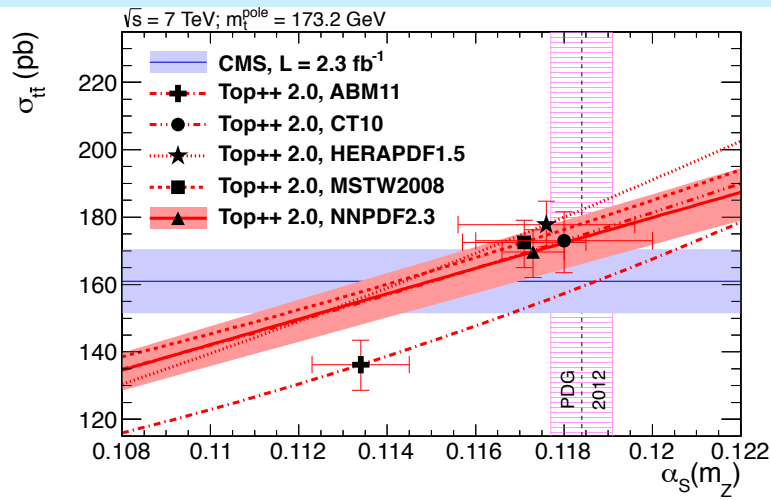
Top-quark pairs at the LHC are produced predominantly in gluon-gluon fusion  
The cross section  $\sigma_{t\bar{t}}$  depends on the value of  $m_t$ ,  $\alpha_s$ , and the gluon distribution,  $g(x)$

Top-pair cross-section as a function of  $m_t$



Dependence  $\sigma_{t\bar{t}}^{\text{exp}}(m_t)$  assumes  $m_t^{\text{MC}} \equiv m_t^{\text{pole}}$

Top-pair cross-section as a function of  $\alpha_s(M_Z)$



Dependence  $\sigma_{t\bar{t}}^{\text{exp}}(\alpha_s)$  found to be small

### CMS Measurement [1]

$\sigma_{t\bar{t}}^{\text{exp}} = 161.9 \pm 6.7$  pb;  
 $m_t^{\text{MC}} = 172.5$  GeV,  
 $\alpha_s(M_Z) = 0.1180$   
Additional uncertainty of 1 GeV [6]  
due to assumption  $m_t^{\text{MC}} \equiv m_t^{\text{pole}}$

### NNLO + NNLL prediction [2,3]

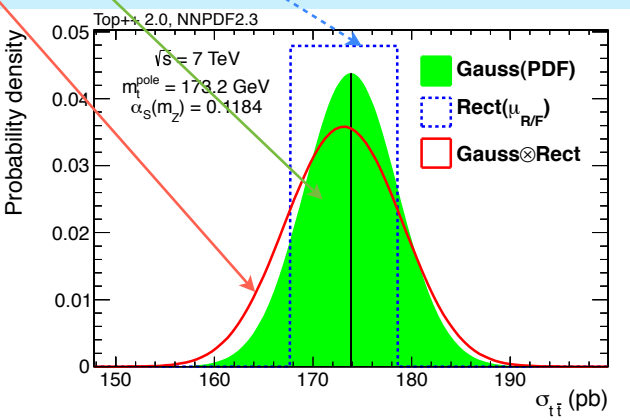
contains  $q\bar{q}, qq', q\bar{q}', qq \rightarrow t\bar{t}+X$   
 $gg \rightarrow t\bar{t}+X, gg \rightarrow t\bar{t}+X$   
QCD scales:  $\mu_r = \mu_f = m_t$   
scale variation  $0.5\mu < \mu_r, \mu_f < 2\mu$   
5 sets of NNLO PDFs used

Both  $m_t$  and  $\alpha_s$  alter the  $\sigma_{t\bar{t}}$  prediction such that any variation of one parameter can be compensated by a variation of the other. Therefore simultaneous extraction of  $\alpha_s$  and  $m_t$  is not possible. For the determination of  $m_t$ ,  $\alpha_s$  is fixed to the world average,  $\alpha_s(M_Z) = 0.1184 \pm 0.0007$  [4]. For the determination of  $\alpha_s$ ,  $m_t$  is fixed to the latest Tevatron average,  $m_t = 173.18 \pm 0.94$  GeV [5] with an additional uncertainty of 1 GeV [6] due to the assumption  $m_t^{\text{MC}} \equiv m_t^{\text{pole}}$ .

Probabilistic approach: maximum of marginalized posterior  $P(x) = \int f_{\text{exp}}(\sigma_{t\bar{t}}|x) f_{\text{th}}(\sigma_{t\bar{t}}|x) d\sigma_{t\bar{t}}$ ,  $x = m_t$  or  $\alpha_s(M_Z)$

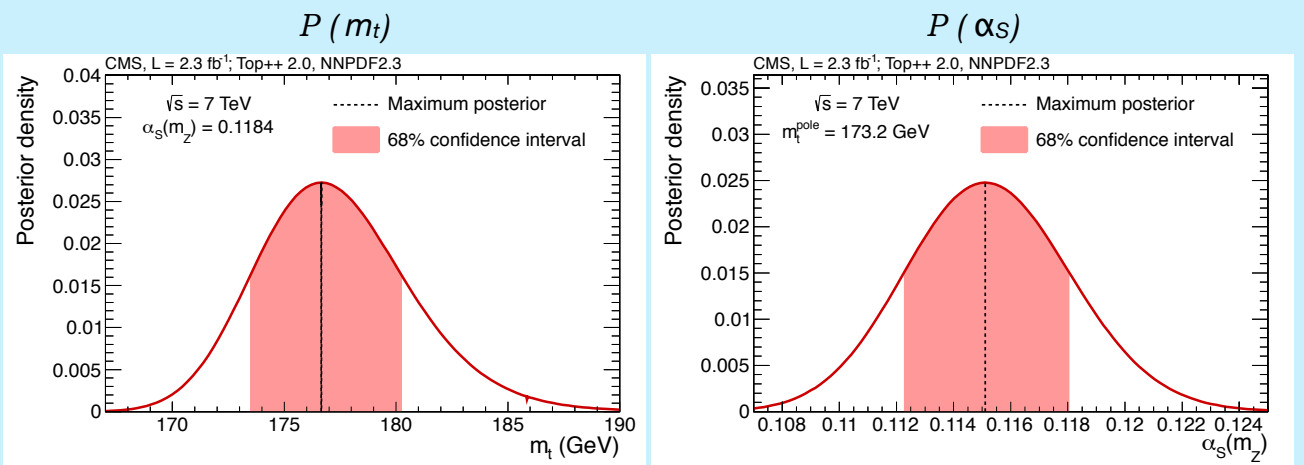
Probability function for predicted cross section

$$f_{\text{th}}(\sigma_{t\bar{t}}) = \mathcal{G}(\delta_{\text{PDF}}) \otimes \text{rect}(\sigma_{t\bar{t}}|\sigma_{t\bar{t}}^{(l)}, \sigma_{t\bar{t}}^{(h)}) = \frac{1}{2(\sigma_{t\bar{t}}^{(h)} - \sigma_{t\bar{t}}^{(l)})} \left( \text{erf} \left[ \frac{\sigma_{t\bar{t}}^{(h)} - \sigma_{t\bar{t}}}{\sqrt{2} \delta_{\text{PDF}}} \right] - \text{erf} \left[ \frac{\sigma_{t\bar{t}}^{(l)} - \sigma_{t\bar{t}}}{\sqrt{2} \delta_{\text{PDF}}} \right] \right)$$



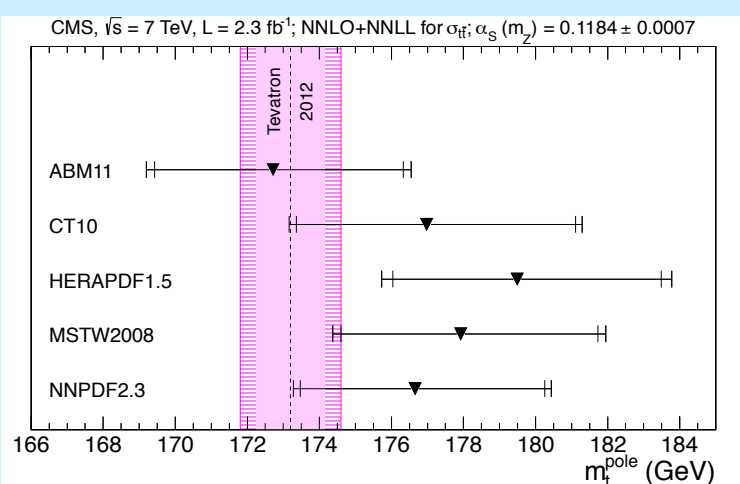
Experimental measurement is represented by Gaussian probability function  $f_{\text{exp}}(\sigma_{t\bar{t}})$

Most probable  $m_t$  or  $\alpha_s(M_Z)$  are obtained from maximum of marginalized posterior:

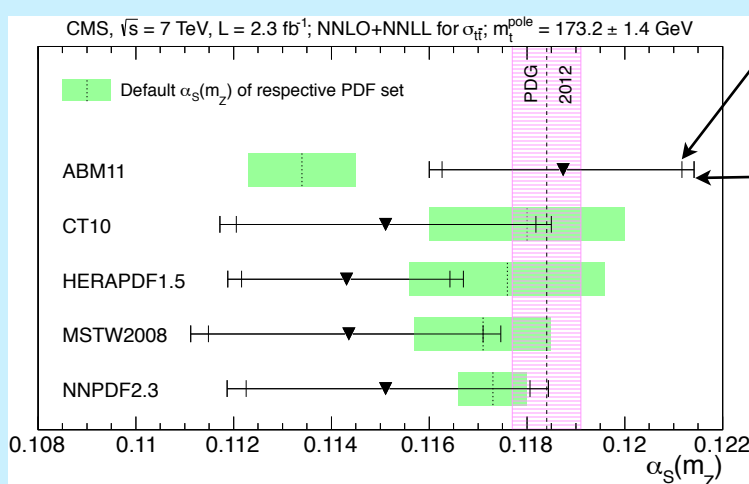


For each PDF set, the most probable values of  $m_t$  and  $\alpha_s(M_Z)$  are obtained

Values of  $m_t$  obtained by confronting  $\sigma_{t\bar{t}}^{\text{exp}}$  and  $\sigma_{t\bar{t}}^{\text{th}}$



Values of  $\alpha_s(M_Z)$  obtained by confronting  $\sigma_{t\bar{t}}^{\text{exp}}$  and  $\sigma_{t\bar{t}}^{\text{th}}$



inner error bars:  
uncertainty on  $\sigma_{t\bar{t}}^{\text{exp}}$ ,  $E_{\text{beam}}^{\text{LHC}}$ ,  
PDF and scale variation in  $\sigma_{t\bar{t}}^{\text{th}}$

outer error bars: uncertainty on  $m_t$   
and  $\alpha_s(M_Z)$  (world average)

Results agree with world average

and consistent for different PDFs

Theory uncertainty (scales)  $\sim 1\%$

Using NNPDF2.3: pole mass of the top quark  $m_t = 176.7^{+3.8}_{-3.4}$  GeV, strong coupling constant  $\alpha_s(M_Z) = 0.1151^{+0.0033}_{-0.0032}$